

Decoherence in a bosonic Josephson junction

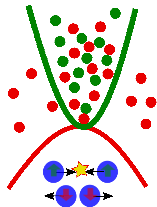
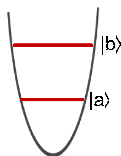
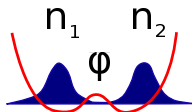
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Quantum Technologies III
10 September 2012, Warsaw

Outline

- Motivation
- N particles in two modes
- 'Typical' evolution - creation of entanglement
- Effect of particle losses



Motivation

'better' statistics = with uncertainties smaller than shot noise

NEEDED:

entangled states

limits due to decoherence and dephasing

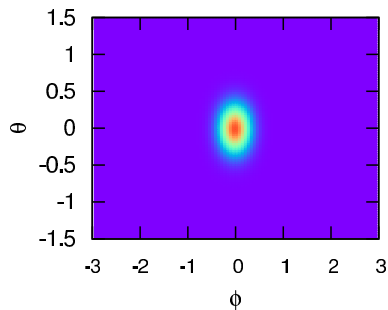
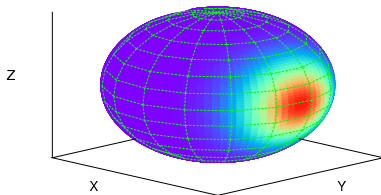
here: effect of particle losses

Initial states

$$|\phi = 0\rangle = \left(\frac{\hat{a}^\dagger + \hat{b}^\dagger}{\sqrt{2}}\right)^N |0\rangle = \sum_{n=0}^N \sqrt{\frac{1}{2^N} \binom{N}{n}} |n, N-n\rangle$$

Initial states

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$$\hat{S}_z = (\hat{n}_a - \hat{n}_b) / 2$$

Evolution

Hamiltonian (**no tunneling**):

$$\begin{aligned}\hat{H} &= \frac{\chi_a}{2} (a^\dagger a^\dagger a a) + \chi_{ab} (a^\dagger a b^\dagger b) + \frac{\chi_b}{2} (b^\dagger b^\dagger b b) = \\ &= \chi \hat{S}_z^2 + f(\hat{N})\end{aligned}$$

If $|\psi\rangle = \sum_{n=0}^N c_n |n, N-n\rangle$, then

$$\begin{aligned}|\psi(t)\rangle &= \sum_{n=0}^N c_n e^{-i(\chi_a - \chi_{ab})(a^\dagger a^\dagger a a)t} \\ &\quad \times e^{-i(\chi_b - \chi_{ab})(b^\dagger b^\dagger b b)t} |n, N-n\rangle\end{aligned}$$

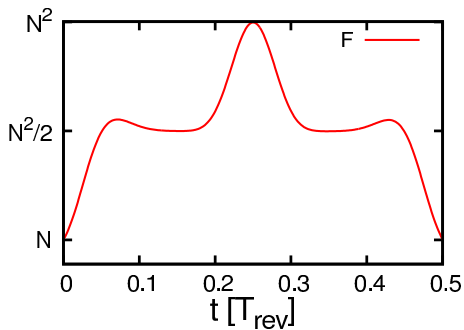
Evolution of a coherent state

Fisher information F

$$F = F(\hat{\rho})$$

$F > N$ - useful
entangled state

$$\max F = N^2$$

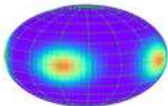


[1] P. Hyllus *et al.* Phys. Rev. A, **82** (2010) 012337

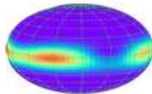
Particle losses

Bloch sphere?

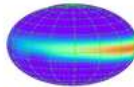
N atoms



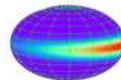
N-1 atoms



N-2 atoms



N-3 atoms



$$\hat{\rho} = p_N \hat{\rho}_N + p_{N-1} \hat{\rho}_{N-1} + p_{N-2} \hat{\rho}_{N-2} + \dots + p_0 \hat{\rho}_0$$

Master equation

$$\underbrace{\partial_t \hat{\rho} = -\frac{i}{\hbar} [\hat{H}, \hat{\rho}]}_{\text{von Neuman equation}} + \underbrace{\mathcal{L}_1 \hat{\rho} + \mathcal{L}_2 \hat{\rho} + \mathcal{L}_3 \hat{\rho}}_{\text{particle losses}},$$

where $\mathcal{L}_n = \mathcal{L}_n^{(a)} + \mathcal{L}_n^{(b)}$

$$\mathcal{L}_n^{(a)} \hat{\rho} = \gamma_n [\hat{a}^n, \hat{\rho} (\hat{a}^\dagger)^n] + \gamma_n [\hat{a}^n, \hat{\rho} (\hat{a}^\dagger)^n]$$

Exact solution

Master equation

$$\underbrace{\partial_t \hat{\rho} = -\frac{i}{\hbar} [\hat{H}, \hat{\rho}]}_{\text{von Neuman equation}} + \underbrace{\mathcal{L}_1 \hat{\rho} + \mathcal{L}_2 \hat{\rho} + \mathcal{L}_3 \hat{\rho}}_{\text{particle losses}},$$

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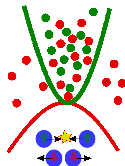
$$\mathcal{L}_n^{(a)} \hat{\rho} = \gamma_n [\hat{a}^n, \hat{\rho} (\hat{a}^\dagger)^n] + \gamma_n [\hat{a}^n, \hat{\rho} (\hat{a}^\dagger)^n]$$

Exact solution

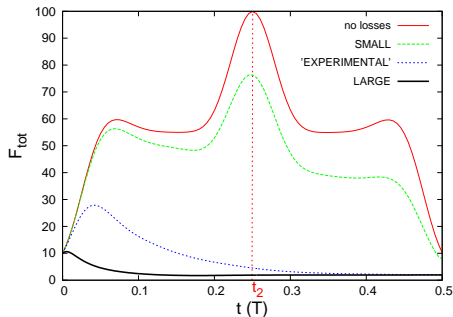
Here: only 2-body losses

γ_1, γ_2 - rates of 2-body losses in "a" and "b" mode

mode



Fisher information with losses



From macroscopic to mesoscopic!

Fisher information with losses

number of lost atoms $\sim 20\%$

γ_1, γ_2 - loss rates in the first and the second mode

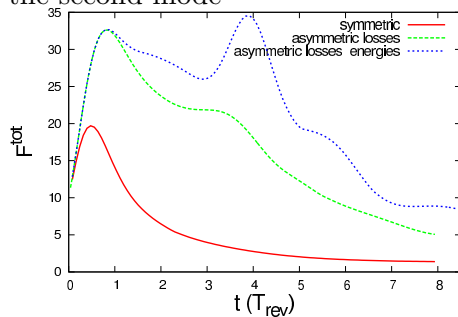
χ_1, χ_2 - interaction energy of a pair of atoms in the first and the second mode

Fisher information with losses

number of lost atoms $\sim 20\%$

γ_1, γ_2 - loss rates in the first and the second mode

χ_1, χ_2 - interaction energy of a pair of atoms in the first and the second mode



Subspace with N atoms

Case, when no losses occurred
(although they were possible)

$p(N, t)$ - probability of an loss event at time t in the cloud with initially N atoms

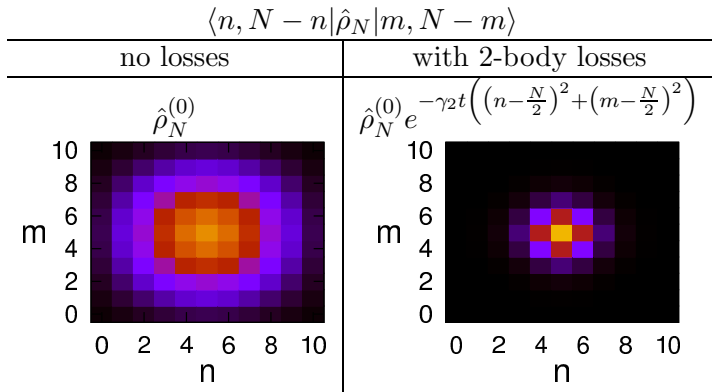
2-body losses

IF $N_1 > N_2$ THEN

$$p(N_1, t) > p(N_2, t)$$

$|0, N\rangle$ less stable, than $|N/2, N/2\rangle$

Do nothing and gain !



Final state = $\left| \frac{N}{2}, \frac{N}{2} \right\rangle$ (but with probability $p_N \ll 1$)

process nr 1: 'Gaussian shrinking'

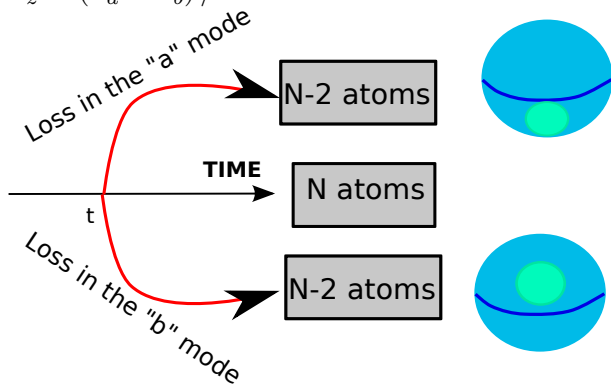
Subspace with $N - 2$ atoms

Two body losses only

What happens when **SINGLE** lost event occurred

Subspace with $N - 2$ atoms - after one loss event

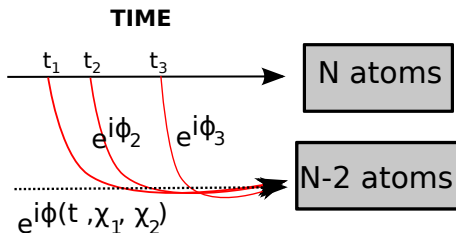
$$\hat{S}_z = (\hat{n}_a - \hat{n}_b) / 2$$



process nr 2: Channeling (with very strong destructive interference for the cat state)

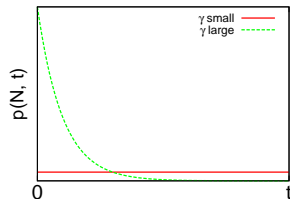
to avoid it: $\gamma_1 \gg \gamma_2$

Subspace with $N - 2$ atoms - after one loss event



state = incoherent mixture of $|\theta_i, \phi_i\rangle$
(estimated for small losses)

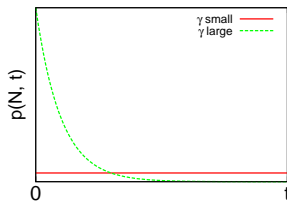
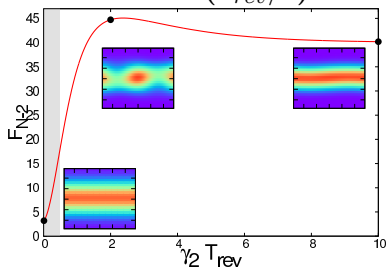
process nr 3: phase noise



$$\Delta\phi \propto \frac{\chi}{2\gamma N}$$

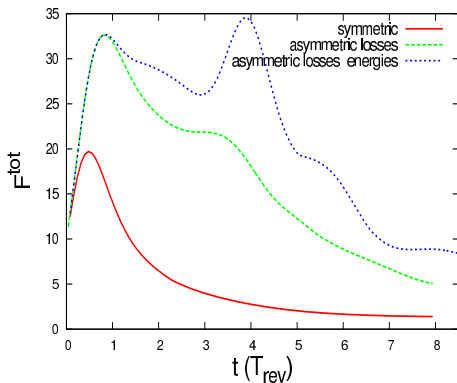
Subspace with $N - 2$ atoms - after one loss event

Fixed time ($T_{rev}/2$) but different loss rates



state = incoherent mixture of $|\theta_i, \phi_i\rangle$ $\Delta\phi \propto \frac{\chi}{2\gamma N}$
 to avoid it: increase γ or suppress χ

Fisher information with losses



$\gamma_1 = 0$ no channeling

$\gamma_1 = 0$ and $\chi_2 = 0$ no phase noise

subspaces with smaller number of atoms

Summary

- ➊ Exact solution of the master equation
- ➋ Decoherence in subspaces with different number of atoms
 - 'Gaussian shrinking'
 - Phase noise
 - Destructive interference (atoms lost from a or b ?)
- ➌ Huge advantage for highly asymmetric losses, using Feshbach resonances
- ➍ Gain via post-selection (?)

TODO-list

- scaling with N
- beyond B-H model
- experimental conditions (phase noise, finite temperature)

Group



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