THE UNITARY GAS:
SYMMETRY PROPERTIES AND APPLICATIONS

Yvan Castin, Félix Werner, Christophe Mora
LKB and LPA, Ecole normale supérieure (Paris, France)
Ludovic Pricoupko
LPTMC, Université Paris 6

GENERAL CONTEXT

The physical system:

- Fermionic atoms with two internal states $\uparrow, \downarrow$
- Short-range interactions between $\uparrow$ and $\downarrow$ controlled by a magnetic Feshbach resonance
- Arbitrary values for the numbers $N_\uparrow$, $N_\downarrow$
- Intense experimental studies (Thomas, Salomon, Jin, Ketterle, Grimm, Hulet, Zwierlein...), e.g. BEC-BCS crossover (Leggett, Nozières, Schmitt-Rink, Sa de Melo,...)

What is not discussed here:

- The actual many-body state of the system: superfluid or normal
- The particularly intriguing strongly polarized case $N_\uparrow \gg N_\downarrow$: Polaronic physics, see talk by C. Trefzger

OUTLINE OF THE TALK

- **What is the unitary gas ?**
- **Simple consequences of scaling invariance**
- **Dynamical consequences: $SO(2,1)$ hidden symmetry in a trap**
- **Separability in hyperspherical coordinates**
- **Does the unitary gas exist ?**
- **First deviations from unitary limit**

WHAT IS THE UNITARY GAS ?

DEFINITION OF THE UNITARY GAS

- **Opposite spin two-body scattering amplitude**

$$f_k = -\frac{1}{ik} \quad \forall k$$

- **"Maximally" interacting: Unitarity of S matrix imposes $|f_k| \leq 1/k$.**
- **In real experiments with magnetic Feshbach resonance:**

$$-\frac{1}{f_k} = \frac{1}{a} + ik - \frac{1}{2}k^2 r_e + O(k^4 b^3)$$

unitary if "infinite" scattering length a and "zero" ranges:

$$k_{\rm typ}|a| > 100, k_{\rm typ}|r_e| \text{ and } k_{\rm typ} b < \frac{1}{100}$$

imposing $|a| > 10$ microns for $r_e \sim b \sim$ a few nm.

- **All these two-body conditions are only necessary.**

THE ZERO-RANGE WIGNER-BETHE-PEIERLS MODEL

- Interactions are replaced by contact conditions.
- For $r_{ij} \to 0$ with fixed ij-centroid $\vec{C}_{ij} = (\vec{r}_i + \vec{r}_j)/2$ different from $\vec{r}_k, k \neq i, j$:

$$\psi(\vec{r}_1, \ldots, \vec{r}_N) = \left(\frac{1}{r_{ij}} - \frac{1}{a}\right) A_{ij}[\vec{C}_{ij}; (\vec{r}_k)_{k\neq i,j}] + O(r_{ij})$$

- Elsewhere, non interacting Schrödinger equation

$$E\psi(\vec{X}) = \left[-\frac{\hbar^2}{2m}\Delta_{\vec{X}} + \frac{1}{2}m\omega^2 X^2\right]\psi(\vec{X})$$

  with $\vec{X} = (\vec{r}_1, \ldots, \vec{r}_N)$.
- Odd exchange symmetry of ψ for same-spin fermion positions.
- Unitary gas exists iff Hamiltonian is self-adjoint.

SIMPLE CONSEQUENCES OF SCALING INVARIANCE

SCALING INVARIANCE OF CONTACT CONDITIONS

$$\psi(\vec{X}) \underset{r_{ij}\to 0}{=} \frac{1}{r_{ij}} A_{ij}[\vec{C}_{ij};(\vec{r}_k)_{k\neq i,j}] + O(r_{ij})$$

- Domain of Hamiltonian is scaling invariant: If ψ obeys the contact conditions, so does ψ_λ with

$$\psi_\lambda(\vec{X}) \equiv \frac{1}{\lambda^{3N/2}}\psi(\vec{X}/\lambda)$$

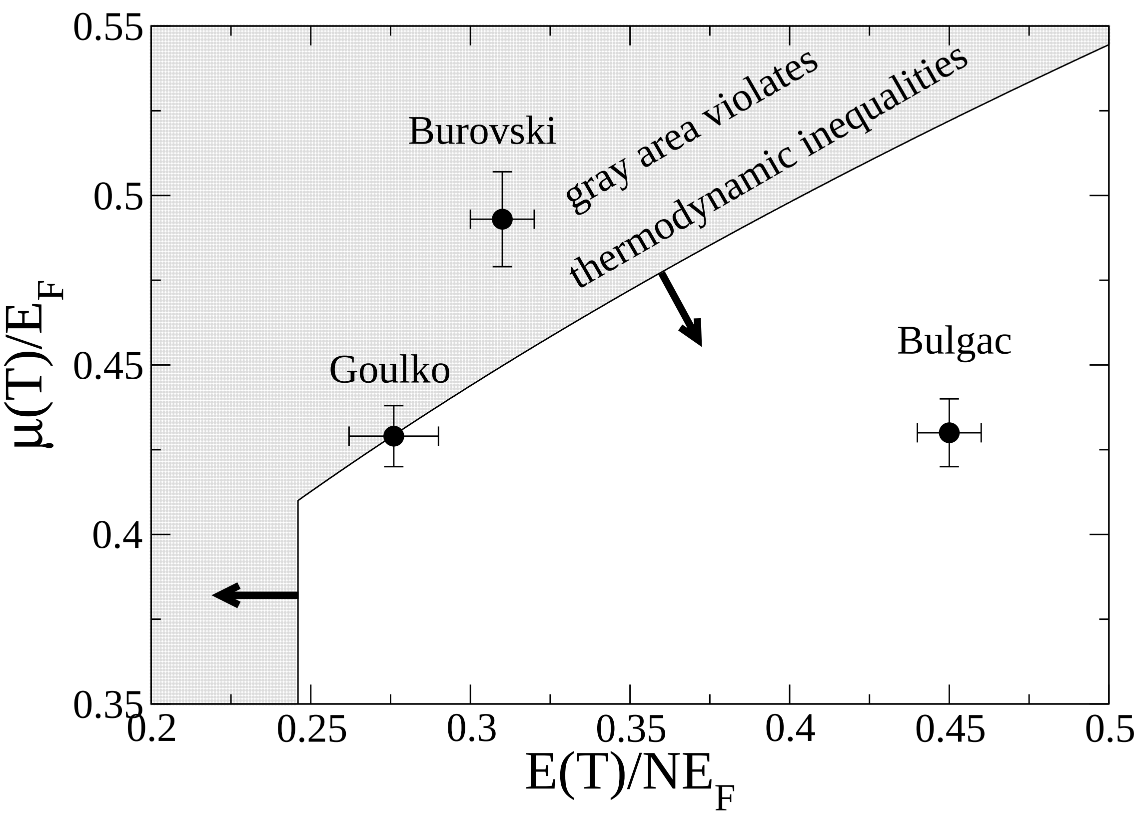
- Consequences (also true for the ideal gas):

free space	box (periodic b.c.)	harm. trap
no bound state$^{(*)}$	$PV = 2E/3$ $^{(**)}$	virial $E = 2E_{\rm harm}$ $^{(***)}$

$^{(*)}$ If ψ of eigenenergy E, ψ_λ of eigenenergy E/λ^2. Square integrable eigenfunctions (after center of mass removal) correspond to point-like spectrum, for selfadjoint H. $^{(**)}$ $E(N, V\lambda^3, S) = E(N, V, S)/\lambda^2$, then take derivative in $\lambda = 1$. $^{(***)}$ For eigenstate ψ, mean energy of ψ_λ stationary in $\lambda = 1$.

TEST FOR QUANTUM MONTE CARLO

For the unpolarized gas in thermodynamic limit, using Carlson's 2009 upper bound on the ground state energy $[\xi = \mu(T=0)/E_F \leq 0,41]$:



DYNAMICAL CONSEQUENCES:

$\boldsymbol{SO(2,1)}$ HIDDEN SYMMETRY IN A TRAP

IN A TIME-DEPENDENT TRAP

- At $t = 0$: static trap $U(\mathrm{r}) = m\omega^2 r^2/2$, system in eigenstate $\psi_0(\vec{X})$ of energy E.
- For $t > 0$, arbitrary time dependence of trap spring constant, $\omega(t)$. Known solution for ideal gas:

$$\psi(\vec{X}, t) = \frac{e^{-i\theta(t)}}{\lambda^{3N/2}(t)} \exp\left[\frac{im\dot{\lambda}}{2\hbar\lambda} X^2\right] \psi_0(\vec{X}/\lambda(t))$$

with $\ddot{\lambda} = \omega^2\lambda^{-3} - \omega^2(t)\lambda$ and $\dot{\theta} = E\lambda^{-2}/\hbar$.

- This is a gauge plus scaling transform.
- The gauge transform also preserves contact conditions:

$$r_i^2 + r_j^2 = 2C_{ij}^2 + \frac{1}{2} r_{ij}^2$$

so solution also applies to unitary gas!

Y. Castin, Comptes Rendus Physique 5, 407 (2004).

IN THE MACROSCOPIC LIMIT

$$\psi(\vec{X},t) = \frac{e^{-i\theta(t)}}{\lambda^{3N/2}} \exp\left[\frac{im\dot{\lambda}}{2\hbar\lambda}X^2\right]\psi_0(\vec{X}/\lambda)$$

density $\rho(\vec{r},t) = \rho_0(\vec{r}/\lambda)/\lambda^3$	velocity field $\vec{v}(\vec{r},t) = \vec{r}\dot{\lambda}/\lambda$
local temp. $T(\vec{r},t) = T/\lambda^2$	pressure $P(\vec{r},t) = P_0(\vec{r}/\lambda)/\lambda^5$
local entropy per particle	$s(\vec{r},t) = s_0(\vec{r}/\lambda)$

This has to solve the hydrodynamic equations for a normal gas. Entropy production equation:

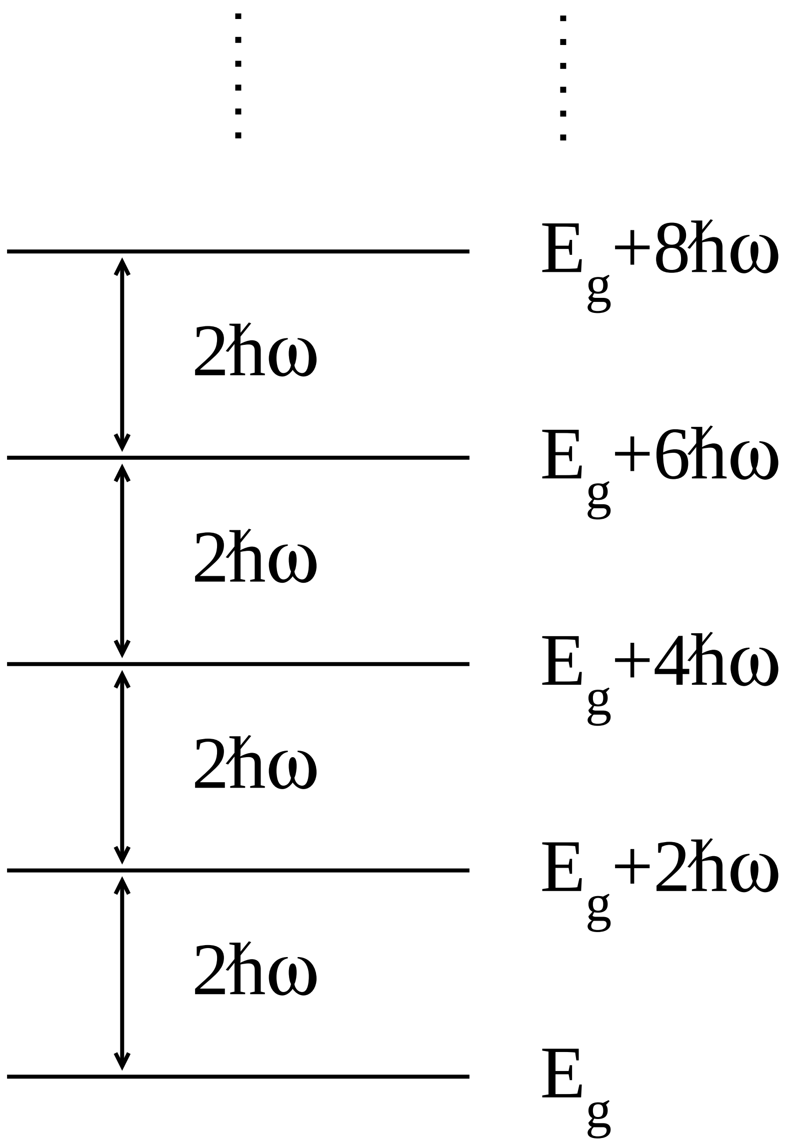
$$\rho k_B T(\partial_t s + \vec{v}\cdot\vec{\nabla}s) = \vec{\nabla}\cdot(\kappa\nabla T) + \zeta(\vec{\nabla}\cdot\vec{v})^2$$
$$+\frac{\eta}{2}\sum_{i,j}\left(\frac{\partial v_i}{\partial x_j} + \frac{\partial v_j}{\partial x_i} - \frac{2}{3}\delta_{ij}\vec{\nabla}\cdot\vec{v}\right)^2$$

so the bulk viscosity is zero: $\zeta(\rho,T) = 0\ \forall T > T_c$. Reproduces the conformal invariance result of Son (2007).

LADDER STRUCTURE OF THE SPECTRUM

- Infinitesimal change of ω for $0 < t < t_f$. For $t > t_f$:
$$\lambda(t) - 1 = \epsilon\, e^{-2i\omega t} + \epsilon^*\, e^{2i\omega t} + O(\epsilon^2)$$
so an udamped mode of frequency 2ω.
- Corresponding wavefunction change:
$$\psi(\vec{X}, t) = \Big[e^{-iEt/\hbar} - \epsilon e^{-i(E+2\hbar\omega)t/\hbar} L_+$$
$$+\epsilon^* e^{-i(E-2\hbar\omega)t/\hbar} L_-\Big]\, \psi_0(\vec{X}) + O(\epsilon^2)$$
- Raising and lowering operators:
$$L_\pm = \pm i\left[\frac{3N}{2i} - i\vec{X}\cdot\partial_{\vec{X}}\right] + \frac{H}{\hbar\omega} - m\omega X^2/\hbar$$
(in red, generator of scaling transform)
- Spectrum=collection of semi-infinite ladders of step $2\hbar\omega$. $SO(2,1)$ hidden symmetry (Pitaevskii, Rosch, 1997).

LADDER STRUCTURE OF THE SPECTRUM (2)



USEFUL MAPPING AND SEPARABILITY

- **Each energy ladder has a ground step of energy E_g, eigenfunction ψ_g.**
- **Integration of $L_-\psi_g = 0$ gives, with $\vec{X} = X\vec{n}$:**

$$\psi_g(\vec{X}\,) = e^{-m\omega X^2/2\hbar} \times \left[X^{E_g/(\hbar\omega)-3N/2} f(\vec{n})\right]$$

- **Limit $\omega \to 0$: mapping to zero energy free space solutions. N.B.: $E_g/(\hbar\omega)$ is a constant.**
- **Free space problem solved for $N = 3$ (Efimov, 1972)... so trapped case also solved (Werner, Castin, 2006).**
- **Also, this is separable in hyperspherical coordinates.**

SEPARABILITY IN HYPERSPHERICAL COORDINATES

SEPARABILITY IN INTERNAL COORDINATES

- **Use Jacobi coordinates to separate center of mass $\vec{C}$**
- **Hyperspherical coordinates:**

$$(\vec{r}_1, \ldots, \vec{r}_N) \leftrightarrow (\vec{C}, R, \vec{\Omega})$$

with $3N-4$ hyperangles $\vec{\Omega}$ and the hyperradius

$$R^2 = \sum_{i=1}^{N} (\vec{r}_i - \vec{C})^2$$

- **Hamiltonian is clearly separable:**

$$H_{\mathrm{internal}} = -\frac{\hbar^2}{2m}\left[\partial_R^2 + \frac{3N-4}{R}\partial_R + \frac{1}{R^2}\Delta_{\vec{\Omega}}\right] + \frac{1}{2}m\omega^2 R^2$$

Do the contact conditions preserve separability ?

- For free space $E = 0$, yes, due to scaling invariance:

$$\psi_{E=0} = R^{s-(3N-5)/2}\phi(\vec{\Omega})$$

$E = 0$ Schrödinger's equation implies

$$\Delta_{\vec{\Omega}}\phi(\vec{\Omega}) = -\left[s^2 - \left(\frac{3N-5}{2}\right)^2\right]\phi(\vec{\Omega})$$

with contact conditions. $s^2 \in$ discrete real set.

- For arbitrary E, Ansatz with $E = 0$ hyperrangular part obeys contact conditions $[R^2 = R^2(r_{ij} = 0) + O(r_{ij}^2)]$:

$$\psi = F(R)R^{-(3N-5)/2}\phi(\vec{\Omega})$$

- Schrödinger's equation for a fictitious particle in 2D:

$$EF(R) = -\frac{\hbar^2}{2m}\Delta_R^{2D}F(R) + \left[\frac{\hbar^2 s^2}{2mR^2} + \frac{1}{2}m\omega^2 R^2\right]F(R)$$

SOLUTION OF HYPERRADIAL EQUATION $(N \geq 3)$

$$EF(R) = -\frac{\hbar^2}{2m}\Delta_R^{2D}F(R) + \left[\frac{\hbar^2 s^2}{2mR^2} + \frac{1}{2}m\omega^2 R^2\right]F(R)$$

- Which boundary condition for $F(R)$ in $R = 0$? Wigner-Bethe-Peierls does not say.
- Key point: particular solutions $F(R) \sim R^{\pm s}$ for $R \to 0$.
- Case $s^2 > 0$: Defining $s > 0$, one discards as usual the divergent solution:

$$F(R) \underset{R\to 0}{\sim} R^s \longrightarrow E_q = E_{\mathrm{CoM}} + (s + 1 + 2q)\hbar\omega, \quad q \in \mathbb{N}$$

- Case $s^2 < 0$: To make the Hamiltonian self-adjoint, one is forced to introduce an extra parameter κ (inverse of a length, calculable via microscopic model). For $s = i|s|$:

$$F(R) \underset{R\to 0}{\sim} (\kappa R)^s - (\kappa R)^{-s}$$

- This breaks scaling invariance of the domain. In free space, a geometric spectrum of N-mers:

$$E_n \propto -\frac{\hbar^2\kappa^2}{m} e^{-2\pi n/|s|}, \quad n \in \mathbb{Z}$$

For $N = 3$, this is the Efimov effect:

- Efimov (1971): Solution for three bosons $(1/a = 0)$. There exists a single purely imaginary $s_3 \simeq i \times 1.00624$.
- Efimov (1973): Solution for three arbitrary particles $(1/a = 0)$. Efimov trimers for two fermions (masse m, same spin state) and one impurity (masse m') iff (Petrov, 2003)

$$\alpha \equiv \frac{m}{m'} > \alpha_c(2;1) \simeq 13.6069$$

DOES THE UNITARY GAS EXIST ?

MINLOS'S THEOREM (1995)

Theorem: *In the $n+1$ fermionic problem, the Wigner-Bethe-Peierls Hamiltonian is self-adjoint and bounded from below iff*

$$(n-1)\frac{2\alpha(1+1/\alpha)^3}{\pi\sqrt{1+2\alpha}}\int_0^{\operatorname{asin}\frac{\alpha}{1+\alpha}} dt\, t\sin t < 1.$$

- α is mass ratio fermion/impurity
- Case $\alpha = 1$: No stable unitary gas for $n > 9$...
- Proof not included in Minlos' paper.
- Proof by Teta, Finco (2010) has a hole.
- A physical test: look for occurrence of $s^2 < 0$ for $n = 3$: four-body Efimov effect !?

ARE THERE EFIMOVIAN TETRAMERS ?

$$E_n^{(4)} \propto -\frac{\hbar^2 \kappa_4^2}{m} e^{-2\pi n/|s_4|} \text{ ?}$$

Negative results for bosons:

- Amado, Greenwood (1973): "There is No Efimov effect for Four or More Particles". Explanation: Case of bosons, there exist trimers, tetramers decay.
- Hammer, Platter (2007), von Stecher, D'Incao, Greene (2009), Deltuva (2010): The four-boson problem (here $1/a = 0$) depends only on κ_3, no κ_4 to add.
- Key point: $N = 3$ Efimov effect breaks separability in hyperspherical coordinates for $N = 4$.

Here, we are dealing with fermions.

OUR DEFINITION OF N-BODY EFIMOV EFFECT

- To find N-body Efimov effect, one simply needs to calculate the exponents s_N, that is to solve the Wigner-Bethe-Peierls model at zero energy:

$$\psi_{E=0}(\vec{r}_1, \ldots, \vec{r}_N) = R^{s_N-(3N-5)/2}\phi(\vec{\Omega})$$

- The N-body Efimov effect takes place iff one of the s_N^2 is < 0.
- This statement makes sense if $\Delta_{\vec{\Omega}}$ self-adjoint for the Wigner-Bethe-Peierls contact conditions: There should be no n-body Efimov effect $\forall n \leq N-1$.

THE $3+1$ FERMIONIC PROBLEM
(Castin, Mora, Pricoupenko, 2010)

- Three fermions (mass m, same spin state) and one impurity (mass m')
- Our def. of 4-body Efimov effect requires a mass ratio

$$\alpha \equiv \frac{m}{m'} < \alpha_c(2;1) \simeq 13.6069$$

- Calculate $E=0$ solution in momentum space. An integral equation for Fourier transform of A_{ij}:

$$0 = \left[\frac{1+2\alpha}{(1+\alpha)^2}(k_1^2+k_2^2) + \frac{2\alpha}{(1+\alpha)^2}\vec{k}_1\cdot\vec{k}_2\right]^{1/2} D(\vec{k}_1,\vec{k}_2)$$

$$+\int \frac{d^3k_3}{2\pi^2}\frac{D(\vec{k}_1,\vec{k}_3)+D(\vec{k}_3,\vec{k}_2)}{k_1^2+k_2^2+k_3^2+\frac{2\alpha}{1+\alpha}(\vec{k}_1\cdot\vec{k}_2+\vec{k}_1\cdot\vec{k}_3+\vec{k}_2\cdot\vec{k}_3)}$$

- D has to obey fermionic symmetry.

RESULTS

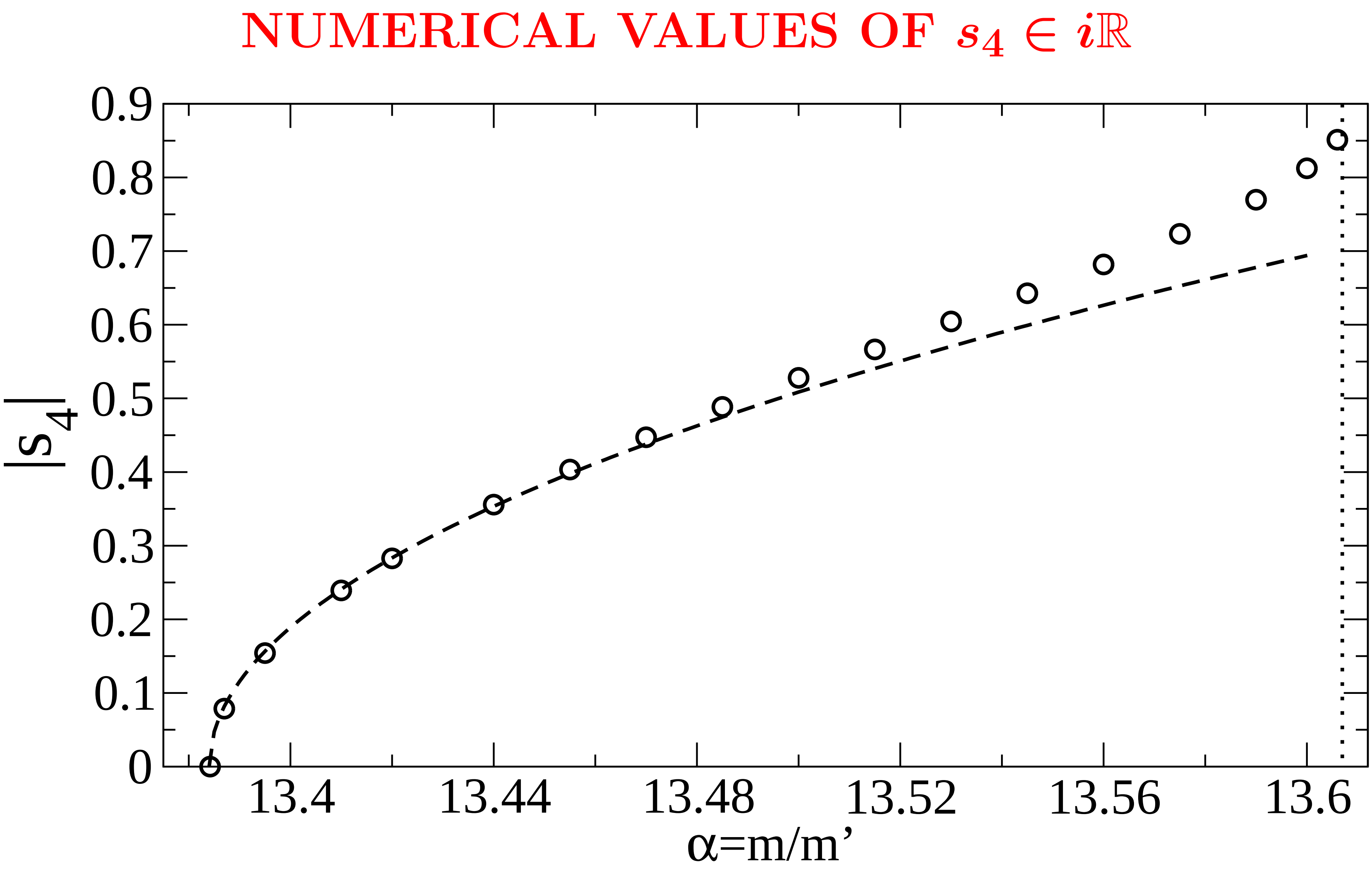
- Four-body Efimov effect obtained for a single s_4, in channel $l = 1$ with even parity:

$$D(\vec{k}_1, \vec{k}_2) = \vec{e}_z \cdot \frac{\vec{k}_1 \times \vec{k}_2}{||\vec{k}_1 \times \vec{k}_2||} \, f_0^{(1)}(k_1, k_2, \theta)$$

in the interval of mass ratio

$$\alpha_c(3;1) \simeq 13.384 < \alpha < \alpha_c(2;1) \simeq 13.607$$

- Strong disagreement with Minlos' critical mass ratio for $n = 3$, $\alpha_c^{\mathrm{Minlos}} \simeq 5.29$
- In experiments: Use optical lattice to tune effective mass of ^{40}K and ^{3}He* away from $\alpha \simeq 13.25$



NUMERICAL VALUES OF $s_4 \in i\mathbb{R}$
0.9
0.8
0.7
0.6
0.5
0.4
0.3
0.2
0.1
0
$|s_4|$
13.4
13.44
13.48
13.52
13.56
13.6
α=m/m'

FIRST DEVIATIONS FROM UNITARITY

FINITE $1/a$ AND FINITE RANGE CORRECTIONS

General relations for the zero-range model:

- Tan relation (generalizing a Lieb relation to 3D):

$$\frac{dE}{d(-1/a)} = \frac{\hbar^2}{4\pi m} \sum_{i<j} \langle A_{ij} | A_{ij} \rangle$$

- The zero-range solution also contains in itself information on finite range corrections (Werner, Castin, 2012):

$$\frac{dE}{dr_e} = 2\pi \sum_{i<j} \langle A_{ij} | \left(H - \frac{p_{ij}^2}{m} \right)_{\vec{r}_i, \vec{r}_j \to \vec{C}_{ij}} | A_{ij} \rangle$$

An experimentally more accessible form:

- Pair distribution function at short distances:

$$\bar{g}^{(2)}_{\uparrow\downarrow}(\vec{r}) = \frac{m}{4\pi\hbar^2} \left[\frac{dE}{d(-1/a)} \left(\frac{1}{r} - \frac{1}{\mathrm{a}} \right)^2 - 2\frac{dE}{dr_e} + O(r) \right]$$

WITHIN A $SO(2,1)$ LADDER

- Reminder of ladder structure:

$$E_q = E_{\mathrm{CoM}} + (s+1+2q)\hbar\omega, \quad q \in \mathbb{N}$$

- N-body problem unsolved: dE/dr_e unknown
- Separability in hyperspherical coordinates leads to explicit expressions (in terms of s and q) for

$$\frac{dE_q/dr_e}{dE_0/dr_e} \quad \text{and} \quad \frac{dE_q/d(1/a)}{dE_0/d(1/a)}$$

- See pioneering work of Moroz (2012).

Large N, unpolarized case:

- Corrections to E_q linear in q: change of breathing frequency 2ω. Agrees with superfluid hydrodynamics (Bulgac, Bertsch)

- **Corrections to E_q quadratic in q: collapse (zero-temperature damping) of breathing mode :**

$$1/t_{\text{collapse}} = \frac{|\delta\omega|}{N^{2/3}} \left| \frac{C_1}{k_F a} + C_2 k_F r_e \right|$$

- **There is a revival of the breathing mode. At half the revival time, a Schrödinger cat state.**